Tailor-made data-driven similarity theories for the temperature structure parameter C_T^2 in the lower atmospheric boundary layer



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What is Optical Turbulence (OT) and why does it matter?

Optical ground-based astronomy and future freespace optical communication (FSOC) suffer from **light getting distorted** when it propagates through the turbulent atmosphere. In astronomy, **turbulent** fluctuations of the atmospheric refractive index, known as optical turbulence (OT), cause blurry images and limit the detection of small objects. FSOC links, which use optical beams to transmit data instead of traditional radio waves, experience reduced data rates or even link interruptions due to laser



Structure function approach used to obtain training data

The structure function coefficient for the indirect method is determined by fitting a linear function to the inertial subrange part of the structure function in log-log space. The 2nd order structure function of sonic temperature is defined as $S_T^2(r) = \langle [T_s(x) - T_s(x+r)]^2 \rangle.$

Assuming frozen turbulence, the spatial increments are obtained from a time signal as $\Delta r \approx \overline{M}/f_s$.



 $\overline{M} \approx 4.6 \,\mathrm{m \, s^{-1}}, \, f_s = 10 \,\mathrm{Hz}, \, \Delta r \approx 0.46 \,\mathrm{m}$



beams getting distorted by OT (see Fig. 1).

Measuring OT needs expensive specialized instruments and simulating OT numerically is expensive due to high required spatial resolution to resolve the inertial subrange of turbulence. That is why, turbulence parameterizations, which estimate turbulent parameters from more readily accessible variables, have been actively researched for decades.

15:00 18:00 21:00 0:00 3:00 6:00 9:00 12:00 15:00 18:00 Time, UTC

Figure 1: Effect of different strengths of optical turbulence on laser beam propagating through the atmosphere.

This poster presents how the **data-driven similarity** theory framework Π-ML [PSB23] is used to obtain physically consistent non-dimensional scalings and parameterizations of OT using only the variables available in specific use cases.

Direct and indirect measurements of optical turbulence strength C_n^2

As basis of our parameterizations, observed C_n^2 data is needed. C_n^2 can be obtained directly or indirectly, where indirect observations from eddy-covariance systems are used in the present study.



Direct: Scintillometer

Measures scintillation, i.e., fluctuations of the received laser intensity (path-averaged)







Indirect: Eddy-covariance system with

ture function coefficients in inertial subrange.

 $C_n^2 = \frac{A_T^2}{\overline{\overline{m}}^2} C_T^2 + \frac{A_T A_q}{\overline{\overline{T}}\overline{\overline{a}}} C_T^2 + \frac{A_q^2}{\overline{\overline{a}}^2} C_q^2$

Obtaining data-driven parameterizations using **II-ML** [PSB23]

1. Collect relevant variables and their dimension (a) traditional: gradients, fluxes, variances (5 min bins)

		0	,	,		`		
	Input features (\boldsymbol{X})							Target (y)
	$\partial \overline{u}_i / \partial z$	$\partial \overline{ heta} / \partial z$	$(u_*)_L$	$\left(\overline{w'\theta'} ight)_L$	$\left(\overline{u_i'^2}\right)_L$	$\left(\overline{\theta'^2}\right)_L$	\mathbf{Z}	C_T^2
Κ		1		1		2		2
L		—1	1	1	2		1	-2/3

3. Train Gradient Boosting Machines (GBMs) for each П-set

Overlapping histograms after scaling indicate good collapse of data in non-dimensional space:



4. Evaluate each model in Π space and normal space

Temporal extrapolation



Traditional parameterizations

Monin-Obukhov similarity theory based [W71]... ... using fluxes: $C_T^2 = (-\overline{w'\theta'}/u_*)^2 z^{-2/3} g(\zeta)$ with stability-dependent similarity function $g(\zeta)$, where $\zeta = z/L$ and $L = -u_*^3 \overline{T}/(\kappa g \overline{w'\theta'})$.

For an in-depth intercomparison, see [PHSB24]



... using gradients: $C_T^2 = \Gamma^2 z^{4/3} f(\operatorname{Ri}_q)$

with stability-dependent similarity function $f(\operatorname{Ri}_q)$

MOST assumptions inherited \rightarrow constant-flux layer!

Physics-based using variances [HB15]: $C_T^2 = 3.2\chi\epsilon^{-1/3}$ with $\epsilon = \frac{(2e)^{3/2}}{B_1 L_M}$, $\chi = \frac{(2e)^{1/2}}{B_2 L_M} \sigma_{\theta}^2$, and master length scale [MY82, NN09, O19] L_M is difficult to $L_M = \left(\frac{1}{L_S} + \frac{1}{L_T} + \frac{1}{L_B}\right)^{-1}$ determine experimentally

Conclusions

- П-ML is a powerful tool to derive non-dimensional scalings for **flow processes** using only the variables available \rightarrow tailor-made
- First, expand complexity of the fitting problem through generation of large number of physically motivated Π -groups and Π -sets.
- Using observational data, machine learning (ML) regression models are fit to assess the capability of each Π-set to parameterize the target.
- The Π-sets and corresponding fitted models are analyzed to **keep** only moderately complex and well-performing combinations.
- For C_T^2 parameterizations, simple enough Π spaces are found to replace ML regression by a simple linear model, ultimately, bringing the complexity back down to an "analytical" level.

Acknowledgments

Results: tailor-made / use case-dependent parameterizations

Traditional: flux and variance-based parameterization



Use case: Outputs of numerical simulations with access to local fluxes and variances, e.g., large eddy simulations or mesoscale simulations with high-order closure.

Variables: local fluxes and variances

Result:
$$\Pi_x = \frac{\sigma_{\theta,L}^2}{\theta_{*,L}^2} \quad \Pi_y = \frac{C_T^2 z^{2/3}}{\theta_{*,L}^2} \quad \theta_{*,L} = \frac{-(\overline{w'\theta'})_L}{u_{*,L}}$$

- Single flux- and variance-based Π expression dominates and is sufficient to collapse data reasonably well (cf. Fig 2).
- Groups containing gradients are discarded by GBM model.
- Scatter may point at missing length scale.
- Scaling is similar to W71 and consistent with theory ([BH22]:
- 1:1 line 60m: R²=0.80, bias=-0.03, cRMSE=0.0 10^{-} ved C_T^2 , $\mathsf{K}^2 \mathsf{m}^$ qo 10^{-} 10^{-2} 10^{-1} 10^{-5} estimated C_T^2 , K² m^{-2/3}

Figure 2: Demonstration of data collapsing in Π space for all levels with simple fitted linear model.

Practical: 1 Hz "temperature only" parameterization



Figure 4: Demonstration of data collapsing in Π space for all levels with simple fitted linear model.

 $C_T^2 \sim \sigma_{\theta}^2 L^{-2/3}$, with $L \sim z$ in neutral conditions). Π groups containing more complex length scales are currently being investigated.

Figure 3: Vertical extrapolation performance of GBM models in log-space using flux-based Π set.

Use case: 1 Hz thermocouples during **field campaign** without wind measurements.

Variables: statistics of (downsampled) 1 Hz temperature time series $\tilde{T}(t)$ and temperature increment signal $\Delta_{\tau}\tilde{T}_s = \tilde{T}_s(t) - \tilde{T}_s(t+\tau)$ lagged by τ samples. Note that C_T^2 for training is still estimated from the 10 Hz signal.

Motivation: No expensive 3D sonic anemometer needed. No wind speed observations needed.

Result:
$$\Pi_x = \frac{\operatorname{var}[\Delta_1 \tilde{T}_s]}{\left(\operatorname{mean}[\Delta_1 \tilde{T}_s]\right)^2} \quad \Pi_y = \frac{C_T^2 z^{2/3}}{\left(\operatorname{mean}[\Delta_1 \tilde{T}_s]\right)^2}$$

- Single Π expression similar to traditional parameterization is sufficient to collapses data very well (cf. Fig. 4).
- Trained GBM models exhibit good but biased vertical extrapolation performance (cf. Fig. 5).



estimated C_T^2 , K² m^{-2/3}

Figure 5: Vertical extrapolation performance of GBM models in log-space using temperature-only Π set.

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